

Three-Step Reversible Cycle

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1 Introduction

In a posting to the CHEMED-L chemical education listserv on 1 August 2003, Pentcho Valev claims to have discovered an exception to the Clausius inequality.

Pentcho's example is that of a closed system containing liquid and gaseous phases of a pure substance in phase equilibrium. The system undergoes a cycle of three reversible steps:

1. adiabatic expansion resulting in vaporization and cooling;
2. heating at constant volume until the initial temperature is regained;
3. compression at constant temperature and pressure until the initial state is reached.

Pentcho claims that if the molar enthalpy of vaporization is large, the integral of dq/T for the cycle can be positive; whereas the Clausius inequality predicts the integral must be zero (because for reversible steps it equals the net entropy change).

I carried out calculations based on high-accuracy experimental data for water available from the U.S. National Institute of Standards and Technology (NIST). These calculations showed that the cyclic integral $\oint dq/T$ for Pentcho's cycle is *zero*, contrary to his claim and in agreement with the Clausius inequality for a cycle of reversible steps.

2 Symbols

n	total amount (moles) of H ₂ O – a constant
n_l	amount (moles) of liquid
n_g	amount (moles) of gas
q, w	heat, work
T, p, V	temperature, pressure, volume
U, H	(internal) energy, enthalpy
$V_m(l), V_m(g)$	molar volumes of liquid and gas
$H_m(l), H_m(g)$	molar enthalpies of liquid and gas
$\Delta_{\text{vap}} V_m$	molar volume of vaporization
$\Delta_{\text{vap}} H_m$	molar enthalpy of vaporization
F	dp/dT along the liquid-gas coexistence curve

3 Mathematical Relations

3.1 General relations

The amounts are related by $n = n_l + n_g$, with n constant, giving

$$n_l = n - n_g \quad dn_l = -dn_g$$

Since the changes in the cycle are reversible, the system is regarded as being in an equilibrium state at all times, with uniform T and p and the liquid and gas in phase equilibrium. Then p (equal to the vapor pressure of the liquid) is a function of T only. Infinitesimal changes in p and T along the liquid–gas coexistence curve are related by the Clapeyron equation $dp/dT = F$ where the function F is given by

$$F = \frac{\Delta_{\text{vap}} H_m}{T \Delta_{\text{vap}} V_m} \quad (1)$$

Normally the intensive properties of a pure phase, such as $V_m(l)$, are functions of p and T . When p is a function of T , however, the intensive properties, including $\Delta_{\text{vap}} V_m$ and $\Delta_{\text{vap}} H_m$, are functions of T only.

The system has two independent variables. The total differentials of V and H , with T and n_g as the independent variables, are given by

$$dV = \left(\frac{\partial V}{\partial T} \right)_{n_g} dT + \left(\frac{\partial V}{\partial n_g} \right)_T dn_g$$

$$dH = \left(\frac{\partial H}{\partial T} \right)_{n_g} dT + \left(\frac{\partial H}{\partial n_g} \right)_T dn_g$$

The partial derivatives $(\partial V/\partial n_g)_T$ and $(\partial H/\partial n_g)_T$ are equal to $\Delta_{\text{vap}} V_m$ and $\Delta_{\text{vap}} H_m$, respectively, so the previous equations become

$$dV = \left(\frac{\partial V}{\partial T} \right)_{n_g} dT + \Delta_{\text{vap}} V_m dn_g \quad (2)$$

$$dH = \left(\frac{\partial H}{\partial T} \right)_{n_g} dT + \Delta_{\text{vap}} H_m dn_g \quad (3)$$

3.2 Step 1

Step 1 is adiabatic: $dq = 0$. The first law applied to Step 1 is

$$dU = dq + dw = 0 - p dV$$

The infinitesimal enthalpy change is

$$dH = d(U + pV) = -p dV + p dV + V dp = V dp = VFdT \quad (4)$$

Equating the expressions for dH in Equations 3 and 4 gives

$$\left(\frac{\partial H}{\partial T}\right)_{n_g} dT + \Delta_{\text{vap}} H_m dn_g = VF dT$$

or

$$dn_g = \frac{VF - (\partial H/\partial T)_{n_g} dT}{\Delta_{\text{vap}} H_m} dT \quad (5)$$

Combining Equations 1, 2, and 5 gives

$$dV = \left[\left(\frac{\partial V}{\partial T}\right)_{n_g} + \frac{V}{T} - \frac{1}{FT} \left(\frac{\partial H}{\partial T}\right)_{n_g} \right] dT \quad (6)$$

3.3 Step 2

The volume is constant in Step 2: $dV = 0$. Equation 2 becomes

$$0 = \left(\frac{\partial V}{\partial T}\right)_{n_g} dT + \Delta_{\text{vap}} V_m dn_g$$

or

$$dn_g = -\frac{(\partial V/\partial T)_{n_g} dT}{\Delta_{\text{vap}} V_m} \quad (7)$$

The energy change is $dU = dq + dw = dq + 0$, and the enthalpy change is

$$dH = d(U + pV) = dq + Vdp = dq + VF dT \quad (8)$$

Equating the expressions for dH in Equations 3 and 8, making substitutions from Equations 1 and 7, and rearranging gives

$$dq = \left[\left(\frac{\partial H}{\partial T}\right)_{n_g} - FT \left(\frac{\partial V}{\partial T}\right)_{n_g} - VF \right] dT \quad (9)$$

4 Calculations

The calculations were made with a program written in Pascal, using experimental data for water downloaded from the NIST Chemistry WebBook at webbook.nist.gov. The data used were values of p , $V_m(l)$, $V_m(g)$, $H_m(l)$, and $H_m(g)$, tabulated at a temperature increment of 0.25 K along the liquid-gas coexistence curve.

4.1 Integration

The integrations referred to in the following paragraphs were carried out as numerical integrations over temperature with a temperature increment or decrement of 0.25 K and using the trapezoidal rule. The integrands of the integrals required values of various quantities evaluated as follows.

At a given temperature, the molar volumes and enthalpies of vaporization were calculated from

$$\Delta_{\text{vap}}V_m = V_m(\text{g}) - V_m(\text{l}) \quad \Delta_{\text{vap}}H_m = H_m(\text{g}) - H_m(\text{l})$$

At a given temperature and for a given value of n_g , the volume and enthalpy were found from

$$V = n_1V_m(\text{l}) + n_gV_m(\text{g}) \quad H = n_1H_m(\text{l}) + n_gH_m(\text{g})$$

The partial derivatives $(\partial V/\partial T)_{n_g}$ and $(\partial H/\partial T)_{n_g}$ were evaluated, at a given value of T and of n_g , as the averages of $\Delta V/\Delta T$ and $\Delta H/\Delta T$ in the two adjacent temperature intervals while keeping n_1 and n_g constant.

4.2 Calculations for one cycle

The calculations for a particular three-step cycle were made with specified values of the total amount of water n , the initial amount of gas, and the initial and final temperatures in Step 1.

For Step 1, n_g after each temperature decrement was found by integrating Equation 5. The volume change in Step 1 was found by integrating Equation 6, and the work was calculated as $-\int p dV$ with dV given by Equation 6.

For Step 2, integration of Equation 7 was used to find n_g after each temperature increment. Equation 9 was used to evaluate q and $\int dq/T$.

In Step 3, taking place at constant T and p , the work is $-p\Delta V$. The heat is $q = \Delta n_g \Delta_{\text{vap}}H_m$ (where Δn_g is the negative of the changes of n_g in Steps 1 and 2), and $\int dq/T$ is equal to q/T .

5 Results

The calculated values listed below are for a cycle in which the total amount of water is $n = 10$ mol, the water is initially all in the liquid phase, and the temperature in Step 1 changes from 350.00 K to 325.00 K.

	T (K)	p (kPa)	V (m ³)	n_g (mol)
initially	350.00	41.682	1.8502×10^{-4}	0
after Step 1	325.00	13.531	8.4535×10^{-2}	0.42420
after Step 2	350.00	41.682	8.4535×10^{-2}	1.21903

	q (kJ)	w (kJ)	$\int dq/T$ (J/K)
Step 1	0	-1.8319	0
Step 2	49.177	0	145.32
Step 3	-50.860	3.5159	-145.31

The important thing to note is that, within the precision of the calculation, the cyclic integral $\oint dq/T$ (the sum of the values in the last column of the second table) is *zero*.